

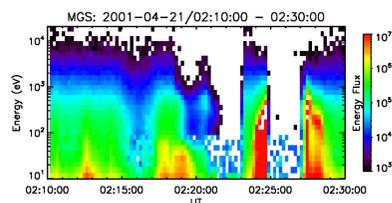
# Electric Field-Driven Currents in the Ionosphere of Mars

## Abstract

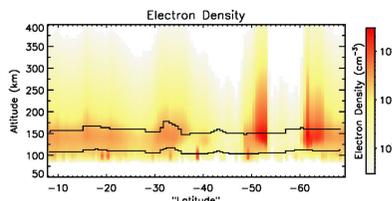
- Previously, we modelled the nightside ionosphere of Mars using precipitating particle data as input to an electron transport model [Fillingim *et al.*, 2007]
- We also calculated horizontal ionospheric currents driven by thermospheric neutral winds [Fillingim *et al.*, 2010]
- Additionally, we estimated the strength of wind-driven electrojets at magnetic cusps created by polarization electric fields formed in the presence of conductivity gradients [Fillingim *et al.*, 2011, in press]
- Here, we extend this previous work by considering ionospheric currents driven by external electric fields
- In the absence of electric field observations, we use the electric field calculated from a global MHD model

## Introduction and Previous Work

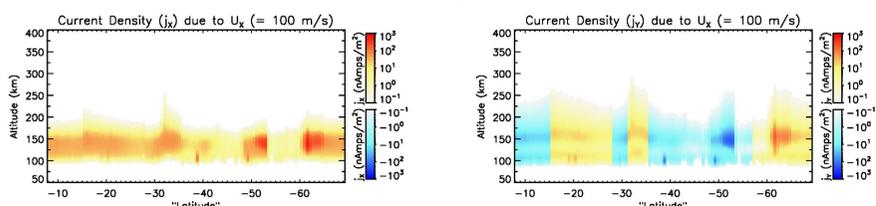
- Start with observed electron energy spectra



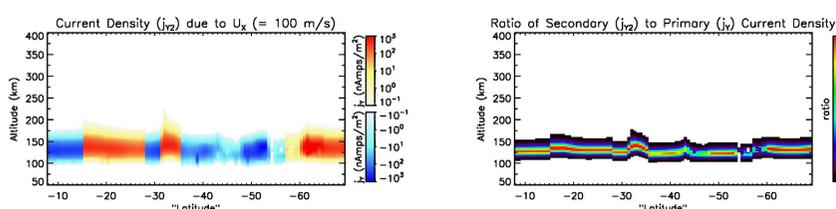
- Calculate ionization rate (Lummerzhim & Lilensten [1994])  
→ neutral atmosphere (MTGCM of Bougher *et al.* [2009])  
→ magnetic field model (Cain *et al.* [2003])
- Compute electron density assuming photochem. equilibrium  
→  $n_e = \sqrt{P(z)/\alpha_{\text{eff}}(z)}$ , where  $\alpha_{\text{eff}}(z)$  is  $\text{O}_2^+$  recombination rate



- Add external force (e.g., neutral wind, electric field)
- Calculate particle velocities,  $\mathbf{v}_{i,e}$ , from equations of motion  
→  $1/n_{i,e} \nabla(n_{i,e} kT_{i,e}) + m_{i,e} \mathbf{g} + q(\mathbf{E} + \mathbf{v}_{i,e} \times \mathbf{B}) - m_{i,e} \mathbf{v}_{in,en}(\mathbf{v}_{i,e} - \mathbf{u}) = 0$
- Calculate currents:  $\mathbf{j} = nq(\mathbf{v}_i - \mathbf{v}_e)$



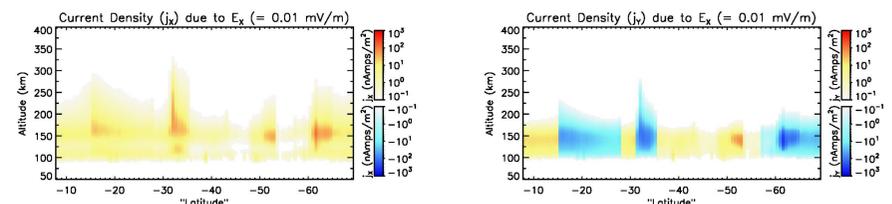
- Assuming  $\mathbf{j}_\perp$  constant across conductivity gradients and  $j_{\parallel} = 0$   
→  $E_{\perp 1} = \Delta j_{\perp 1} / \sigma_P$ ;  $j_{H12} = E_{\perp 1} \sigma_H = \Delta j_{\perp 1} \sigma_H / \sigma_P$ ;  $j_{H12} / j_{\perp 1} \approx \sigma_H^2 / \sigma_P^2$



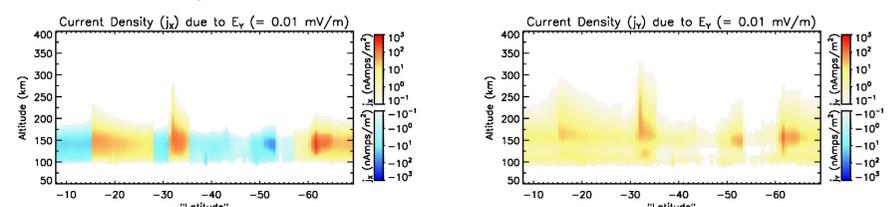
## Methodology and Results

- Perform a similar calculation driven by an externally imposed electric field,  $\mathbf{E}$ , based on Ma *et al.*'s [2004] MHD simulation
- (Grossly over-) Simplifying assumptions:
  - magnetic field is vertical ( $\mathbf{B} = B_z$ ) and constant with altitude
  - magnetic field lines are equipotentials  
→ electric field is horizontal and constant with altitude
  - electric field is constant with latitude on the nightside (actually,  $\mathbf{E}$  varies with latitude, longitude, and time)

- Consider two cases:  $\mathbf{E} = E_x = 0.01$  mV/m (northward field)

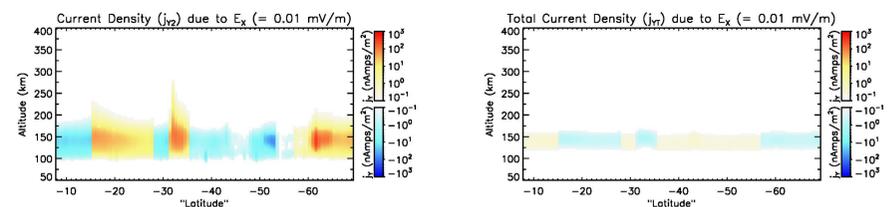


- Case 2:  $\mathbf{E} = E_y = 0.01$  mV/m (westward electric field)

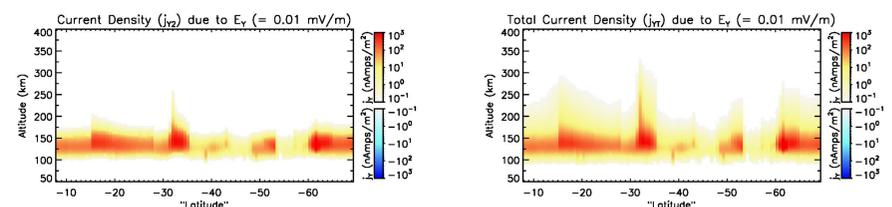


- In both cases, we see two peaks in the Pedersen current as a function of altitude (cf. Opgenoorth *et al.*, 2010) and a change in sign of the Hall current as the polarity of  $\mathbf{B}$  changes
- If we assume  $j_{\parallel} = 0$ , then in steady state  $\nabla \cdot \mathbf{j}_\perp = 0$  →  $dj_x/dx = 0$   
→ gradients in conductivity lead to gradients in  $j_x$   
→ polarization electric fields balance  $j_x$ , induce secondary  $j_{y2}$

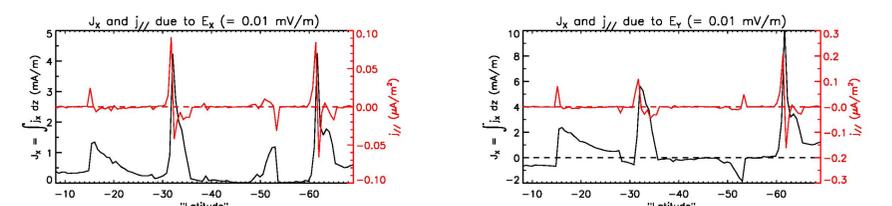
- For Case 1, the secondary  $j_{y2}$  (nearly) cancels the primary  $j_y$



- For Case 2,  $j_{y2}$  strongly reinforces  $j_y$  → "auroral" electrojets



- If we **do not** assume  $j_{\parallel} = 0$ , then  $\nabla \cdot \mathbf{J}_\perp = dj_x/dx = j_{\parallel}$   
where  $J_x$  is the height-integrated current



- Parallel current near cusps → may play a role in acceleration
- All cases have significant Joule heating ( $\mathbf{J} \cdot \mathbf{E} > 0$ ) *except* for Case 1 with polarization electric fields;  $j_{\parallel} = 0$  and  $j_x, E_x \rightarrow 0$